

The $\alpha(v)$ curve *per se* provides a means of estimating the validity of the $\alpha(v)$ approximation. In the limit of low pressures, α should emerge smoothly into the van der Waals b . Hence, as a covolume, α should be about four times greater than the actual volume v_m of the molecules. Hence one should have for the internal energy increase due to the applied pressure p ,

$$E_i = -\frac{1}{4} \int_0^p p d\alpha \cong -\frac{nRT}{4} \int_0^v d\alpha/(v - \alpha).$$

The integral has as an upper limit for CHNO explosives a value of about 2.5 giving an upper limit value for E_i of about 200 cal/g. But Q (the heat of explosion) plus the Hugoniot energy $\frac{1}{2}p(v_1 - v_2)$ is in this upper limit case about 1800 cal g⁻¹ giving $E_i/(Q + \frac{1}{2}p(v_1 - v)) = 0.1$ to 0.15.

One may also estimate E_i directly by the methods used here for thermal expansion and compressibility of solids. Thus

$$E_i = \Delta T = \int_0^p p dv = \int_0^p v_m p \left(\frac{1}{v_m} \frac{dv_m}{dp} \right) dp < v_m \beta p^2 / 2.$$

The theory for molecules of the type considered here also gives $\beta_0 \leq 10^{-6}$ atm⁻¹. Since \bar{v}_m should have an upper limit of 0.5 cm³ g⁻¹, one thus obtains $E_i \sim 200$ cal g⁻¹ in agreement with the above estimate.

According to these estimates therefore, while the $\alpha(v)$ equation of state gives an upper limit for T_2 , the actual value should be no more than 10-15 % (or about 500° K) less. This is, however, in sharp contrast with the approximately 3000° K lower values for T_2 computed for PETN at $\rho_1 = 1.7$ g cm⁻³ by Cottrell and Paterson.¹⁹

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¹ Cook, *J. Chem. Physics*, 1947, **15**, 518; 1948, **16**, 554, 1165; see also Paterson, *J. Chem. Physics*, 1948, **16**, 159, 847.

² Cook, Keyes, Horsley and Filler, *J. Physic. Chem.*, 1954, **58**, 1114.

³ Cook, Horsley, Partridge and Ursenbach, *J. Chem. Physics*, 1956, **24**, 60.

⁴ Cook and Ursenbach, *2nd ONR Symp. Detonation* (Feb. 9-11, 1955, Washington, D.C.), p. 401.

⁵ Paterson, *Research*, 1948, **1**, 221.

⁶ Murgai, *J. Chem. Physics*, 1956, **24**, 635.

⁷ Caldirola, *J. Chem. Physics*, 1946, **14**, 730.

⁸ Gibson, Bowser, Summers, Scott, Cooper and Mason, *2nd ONR Symp. Detonation* (Feb. 9-11, 1955, Washington, D.C.), p. 50.

⁹ Cook, *Bull. no. 53*, **42**, no. 2, 1951, Utah Eng. Expt. Station.

¹⁰ Cook, *J. Chem. Physics*, 1945, **13**, 262; 1946, **14**, 62; *J. Physic. Chem.*, 1947, **51**, 407; *Utah Acad. Sci., Arts Letters*, 1948, **25**, 145.

¹¹ Cook, *Bull. no. 74*, **46**, no. 16, Nov. 30, 1955 (Utah Eng. Expt. Station).

¹² Mott and Gurney, *Electronic Processes in Ionic Crystals* (Oxford, Clarendon Press, 1953).

¹³ Micehalson, *J. Appl. Physics*, 1936, **21**, 536.

¹⁴ Mott and Jones, *Properties of Metals and Alloys* (Oxford, Clarendon Press, 1936).

¹⁵ Seitz, *Modern Theory of Solids* (McGraw Hill Co., 1940).

¹⁶ Dorsey, *Physic. Rev.*, 1907, **25**, 88.

¹⁷ Bridgman, *The Physics of High Pressures* (MacMillan Co., New York, 1951).

¹⁸ Walsh and Christian, *Physic. Rev.*, 1955, **97**, 1544.

¹⁹ Cottrell and Patterson, *Proc. Roy. Soc. A*, 1952, **213**, 214.